



# Modelling hydrogels: building networks in the Mathematical Sciences

Mathematical Interdisciplinary Research at Warwick (MIR@W) Day  
**Warwick Mathematics Institute, 9<sup>th</sup> December 2024**

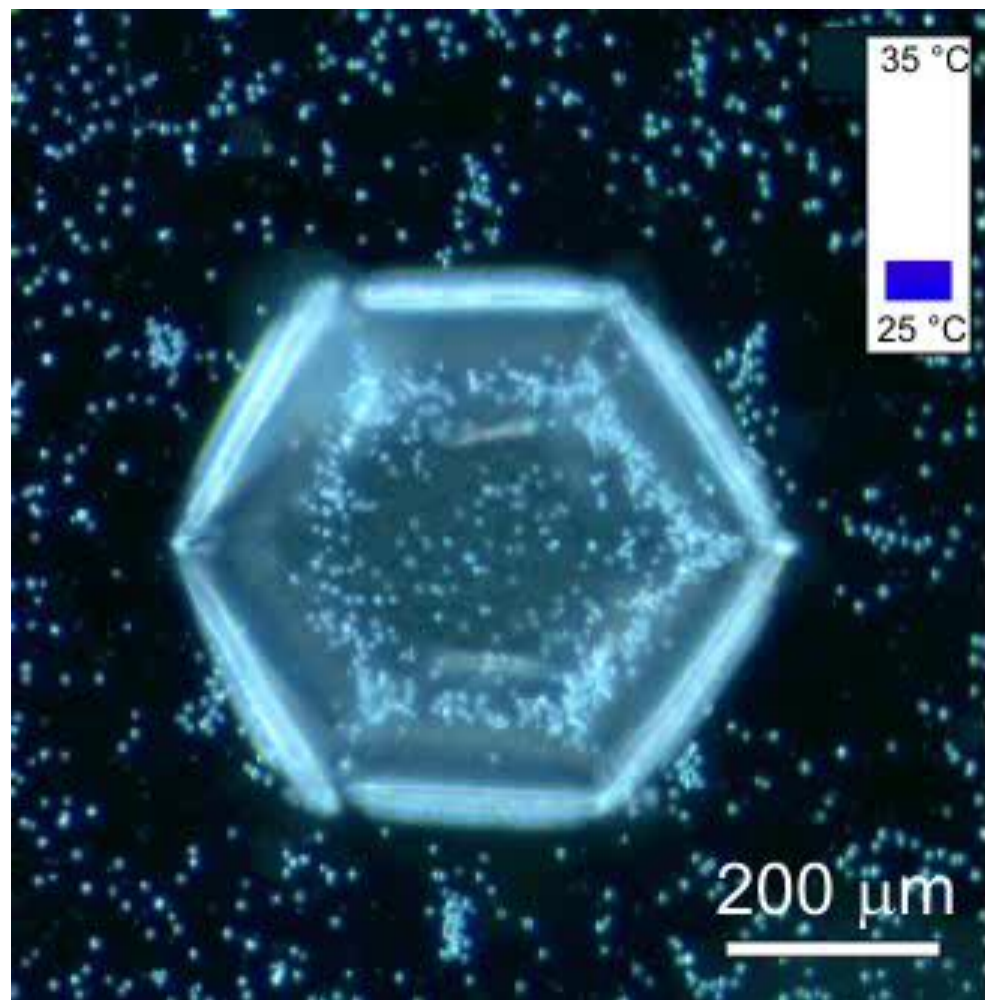
# Deswelling response to temperature changes

building macroscopic models for thermo-responsive gels

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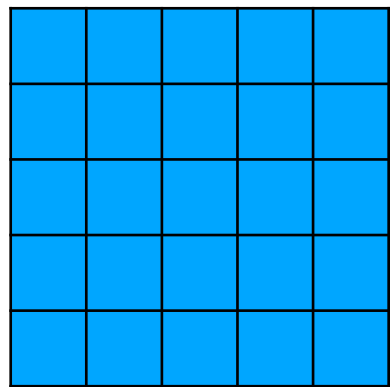


Stoychev *et al.* Soft Matter 7 (2011)

# Thermo-responsive hydrogels

$\phi$  Polymer (volume) fraction  
(1-10% when 'swollen')

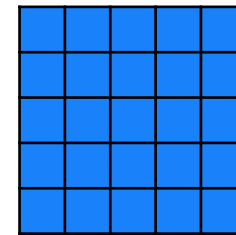
Equilibrium polymer fraction = the degree of swelling reached by a gel placed in water, free to swell, in the absence of mechanical constraints



$$\phi \equiv \phi_0(T_1)$$



$$\phi \equiv \phi_0(T_2)$$



water expelled

- Takes some time – water diffusion is **slow**
- Qualitatively, looks like there's a single swollen state and a single dry state either side of a **lower critical solution temperature (LCST)  $T = T_C$**

...most of the time, anyway, if you ignore Afroze, Nies & Berghmans (2000) or Butler & Montenegro-Johnson (2022)

# Thermo-responsive hydrogels

$\phi$  Polymer (volume) fraction  
(1-10% when 'swollen')

$$\mathcal{W} = \underbrace{\frac{k_B T}{2\Omega_p} [\text{tr}(\mathbf{F}_d \mathbf{F}_d^T) - 3 + 2 \log \phi]}_{\text{neo-Hookean elasticity}} + \underbrace{\frac{k_B T}{\Omega_f} \left[ \frac{1 - \phi}{\phi} \log(1 - \phi) + \chi(\phi, T)(1 - \phi) \right]}_{\text{mixing energy}}$$

- Affinity for water is encoded by the Flory chi parameter measuring the strength of electrostatic attraction between polymer and water

$$\chi(\phi, T) = A_0 + A_1 T + (B_0 + B_1 T)\phi + \mathcal{O}(T^2, \phi^2)$$

- Can fit the parameters here from observations – somewhat crude and very sensitive to experimental error

# Thermo-responsive LENS

$$\mathbf{e}_{ij} = \left[ 1 - (\phi/\phi_{00})^{1/3} \right] \delta_{ij} + \epsilon_{ij} \quad | \quad |\epsilon_{ij}| \ll 1$$

$$\sigma_{ij} = -p\delta_{ij} + \phi \frac{\partial \mathcal{W}}{\partial F_{ik}} F_{jk}$$

Deformation is, at leading order, isotropic, corresponding to swelling or drying.

Gels are characterised by three material parameters dependent on degree of swelling.

see Webber & Worster (2023)

Measure deformation relative to a “swollen” steady state at a temperature well below the LCST, where  $\phi = \phi_{00}$

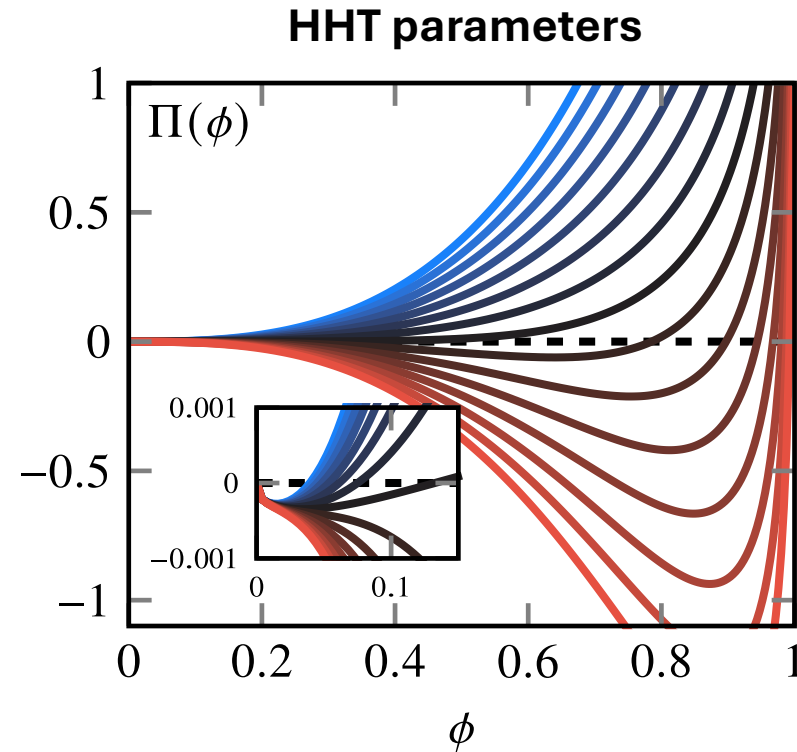
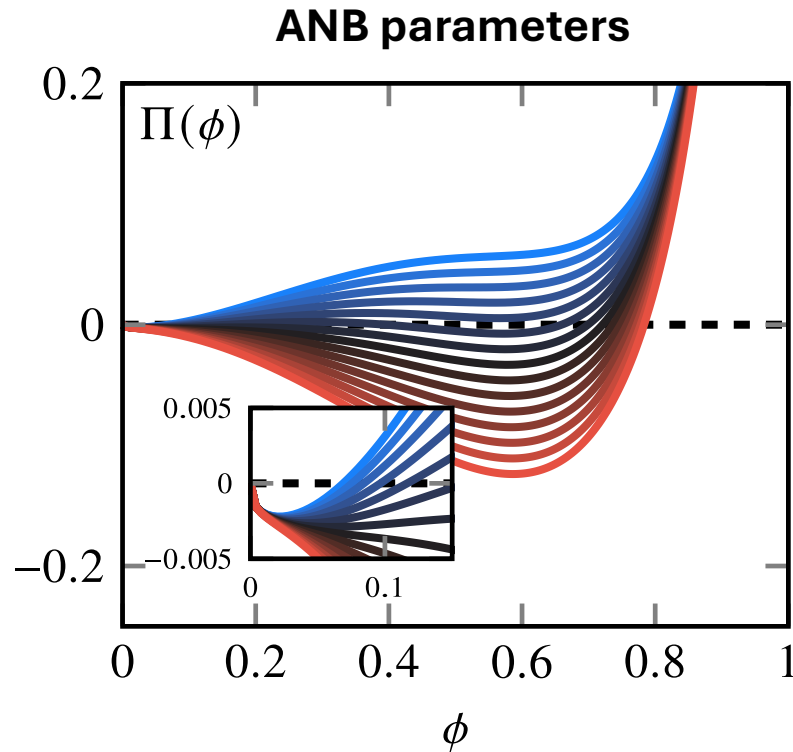
$$\sigma_{ij} = -[p + \Pi(\phi)]\delta_{ij} + 2\mu_s(\phi)\epsilon_{ij}$$

- $p$  is the pervadic/pore/Darcy pressure of the fluid
- $\Pi$  is the **osmotic pressure**
- $\mu_s$  is the **shear modulus**

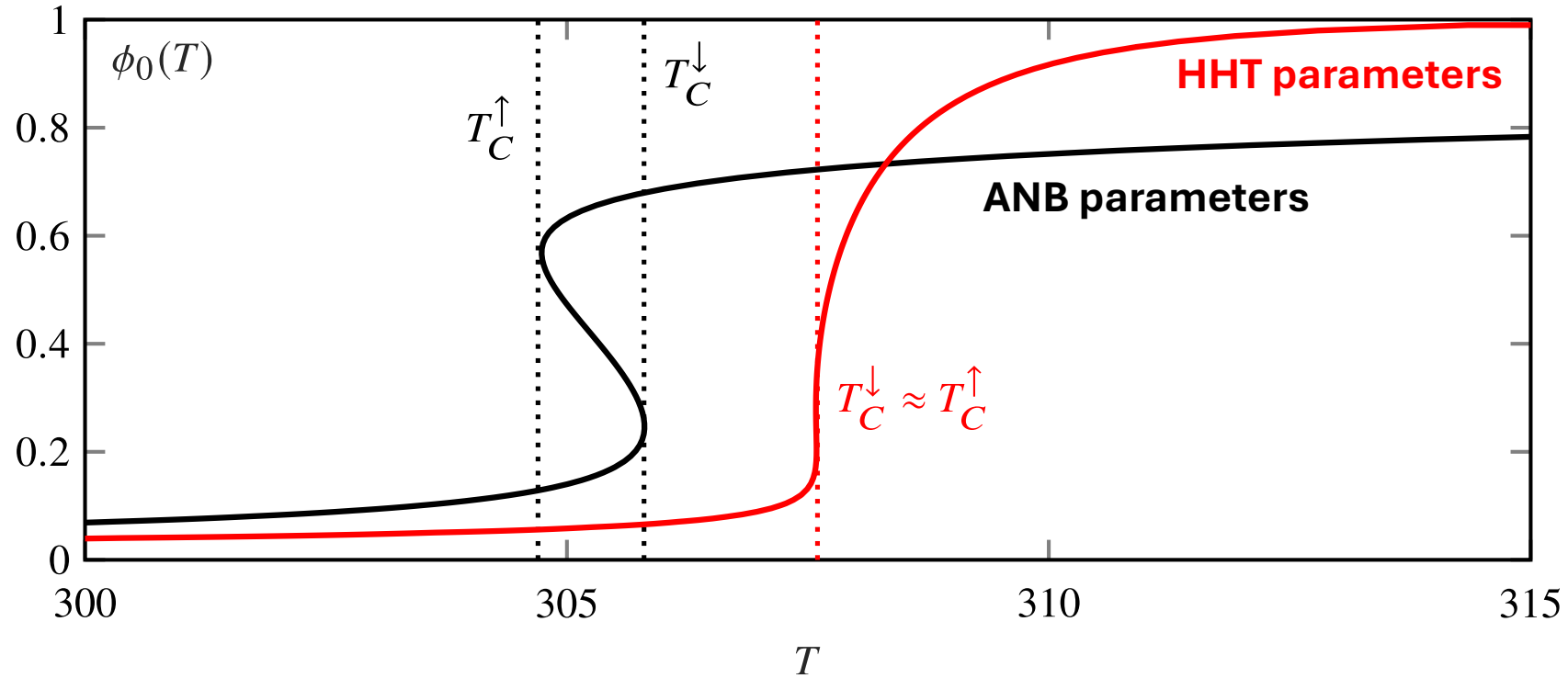
$$\Pi(\phi) = \frac{k_B T}{\Omega_f} \left[ \frac{\Omega_f}{\Omega_p} (\phi - \phi^{1/3}) - \phi - \log(1 - \phi) - \phi^2 \chi + \phi^2 (1 - \phi) \frac{\partial \chi}{\partial \phi} \right]$$
$$\mu_s(\phi) = \frac{k_B T \phi_{00}^{1/3}}{\Omega_p}$$

# Osmotic pressures

$$\sigma_{ij} = -[p + \Pi(\phi)]\delta_{ij} + 2\mu_s(\phi)\epsilon_{ij}$$

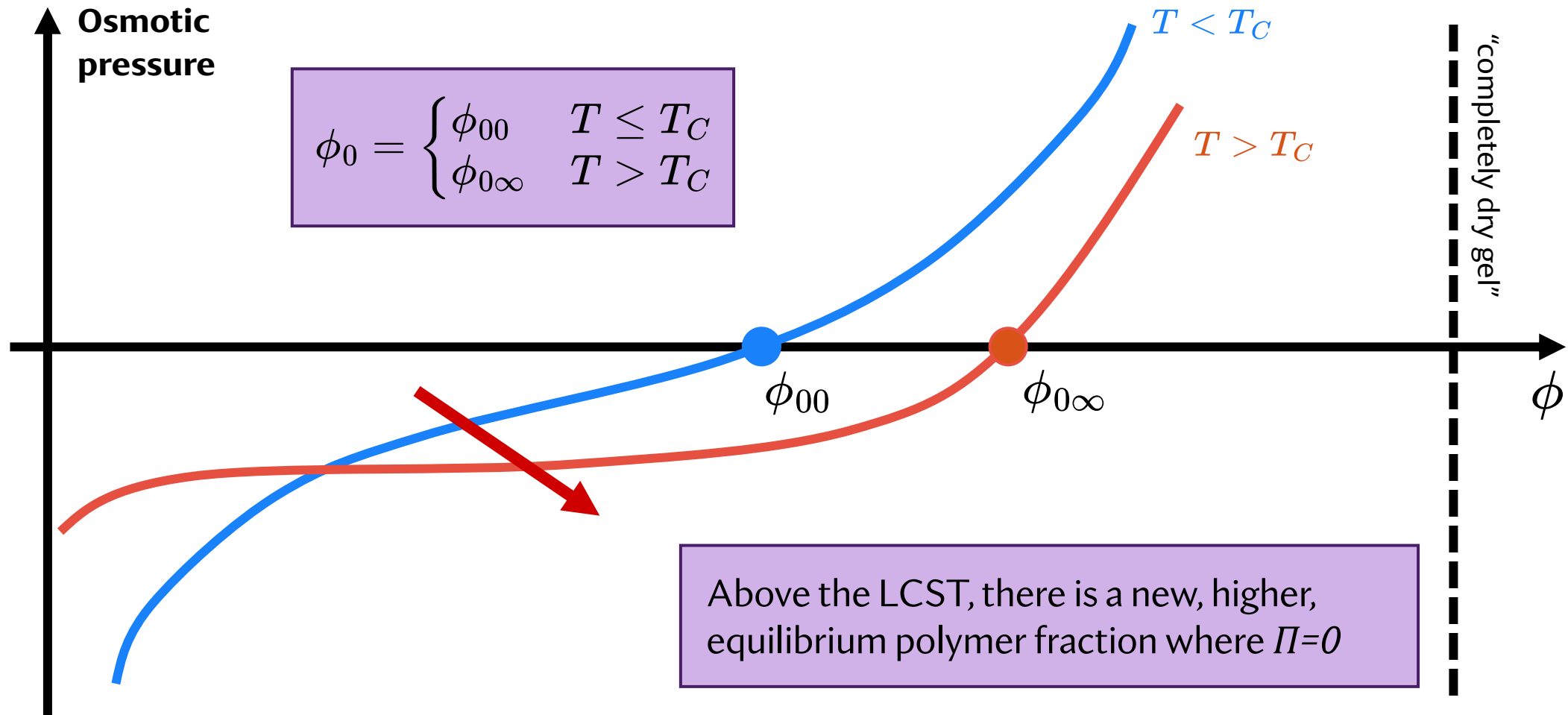


# Osmotic pressures





# Phenomenological observations



# The thermo-responsive LENS framework

## Osmotic components

$$\phi_0(T) = \begin{cases} \phi_{00} & T \leq T_C \\ \phi_{0\infty} & T > T_C \end{cases}$$

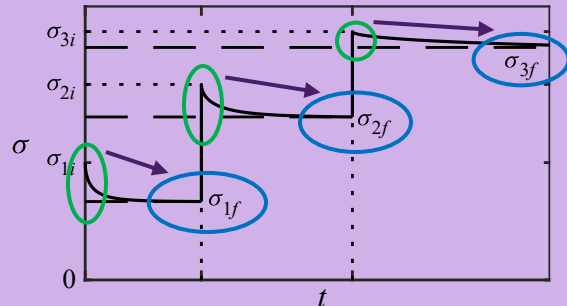
$$\Pi(\phi) = \frac{\phi - \phi_0(T)}{\phi_0(T)} \quad \text{see Doi (2009)}$$

## LENS modelling

$$\boldsymbol{\sigma} = - [p + \Pi(\phi)] \mathbf{I} + 2\mu_s(\phi) \boldsymbol{\epsilon}$$

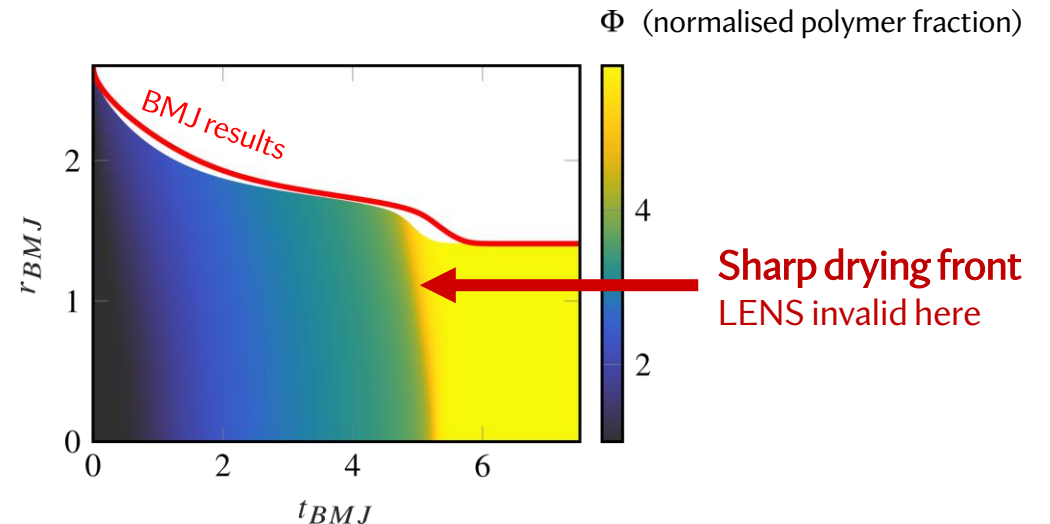
$$\frac{\partial \phi}{\partial t} + \mathbf{q} \cdot \nabla \phi = \nabla \cdot \left\{ \frac{k(\phi)}{\mu_l} \left[ \phi \frac{\partial \Pi}{\partial \phi} + \frac{4\mu_s(\phi)}{3} \left( \frac{\phi}{\phi_0} \right)^{1/3} \right] \nabla \phi \right\}$$

Can feasibly measure osmotic pressure in a macroscopic experiment



**Example:** deswelling hydrogel bead in water

comparison with Butler & Montenegro-Johnson (2022)



# Modelling temperature evolution

- Is this all? Is it reasonable to neglect heat contributions from swelling or drying? How about thermoelasticity?
- Furthermore, how is heat transferred through a responsive gel? Must couple influence from advection, diffusion, heat sources/sinks...

## Heat transfer in porous media

*Treats each phase separately, with a coupling term (e.g.  $H(T_{water} - T_{polymer})$  or similar).*

- Easy physical interpretation
- Matches nicely with continuum models
- Hard to couple heat sources/sinks and quantify coupling between phases

## Thermoelasticity

*Starts from the laws of thermodynamics and the free energy density function from before.*

- Captures potentially complicated thermodynamics
- Hard to interpret some terms: **can LENS help?**

# Modelling temperature evolution

$\mu$  Chemical potential  
 $= \Omega_f p$  (Peppin et al 2005)

Internal energy change  
 Take  $U = cT$

External supply of heat  
 Ignore here

Heat generation (thermoelasticity)

$$\mathbf{F} \approx (\phi/\phi_{00})^{-1/3} \mathbf{I}$$

$$\frac{dU}{dt} = R - \nabla \cdot \mathbf{Q} + \phi^{-1} \sigma \mathbf{F}^{-T} : \frac{d\mathbf{F}}{dt} - \mathbf{J} \cdot \nabla \mu + \mu \frac{dC}{dt}$$

Number density of water molecules

$$C = \frac{1}{\Omega_f} \left( \frac{1}{\phi} - 1 \right)$$

Heat flux

Two components – transport of material *and* Fourier-type conduction

$$\mathbf{Q} = cT \mathbf{q} - c\kappa \nabla T$$

Molecular flux of water

$$\mathbf{J} = \mathbf{u}/\Omega_f = -\frac{k(\phi)}{\mu_l \Omega_f} \nabla p$$

$$c \frac{dT}{dt} = c \mathbf{q} \cdot \nabla T + cT \nabla \cdot \mathbf{q} - \frac{1}{3\phi^2} \frac{\partial \phi}{\partial t} \text{tr} \sigma + \frac{k(\phi)}{\mu_l} |\nabla p|^2 - \frac{p}{\phi^2} \frac{\partial \phi}{\partial t} + c\kappa \nabla^2 T$$

$\text{tr} \sigma = -3P = -3(p + \Pi)$

# Modelling temperature evolution

## Advection

Heat energy carried with deforming gel

$$\frac{D_q T}{Dt} = \frac{\partial T}{\partial t} + \mathbf{q} \cdot \nabla T = \kappa \nabla^2 T + \frac{\Pi(\phi)}{c\phi^2} \frac{\partial \phi}{\partial t} + \frac{k(\phi)}{\mu_l c} |\nabla p|^2$$

## Diffusion

(having assumed spatially-constant diffusivity)

## “Latent heat of swelling or drying”

Energy used up or released in swelling/drying and heat generation by deformation. Dependent on [generalised] osmotic pressure.

## Work done by pore flows

Viscous resistance as water flows through the gel pore spaces

N.B. this reduces down to simple diffusion when the gel isn't reconfiguring

# Modelling temperature evolution

$$\tau = \frac{\mu_l L^2}{k \hat{\Pi}}$$

$$\frac{D_q T}{Dt} = \frac{\partial T}{\partial t} + \mathbf{q} \cdot \nabla T = \kappa \nabla^2 T + \frac{\Pi(\phi)}{c \phi^2} \frac{\partial \phi}{\partial t} + \frac{k(\phi)}{\mu_l c} |\nabla p|^2$$

$\frac{\Delta T}{\tau} \quad \frac{\kappa \Delta T}{L^2} \quad \frac{\hat{\Pi}}{c \tau}$

- Can compare relative importance of heat transfer methods using two non-dimensional numbers,

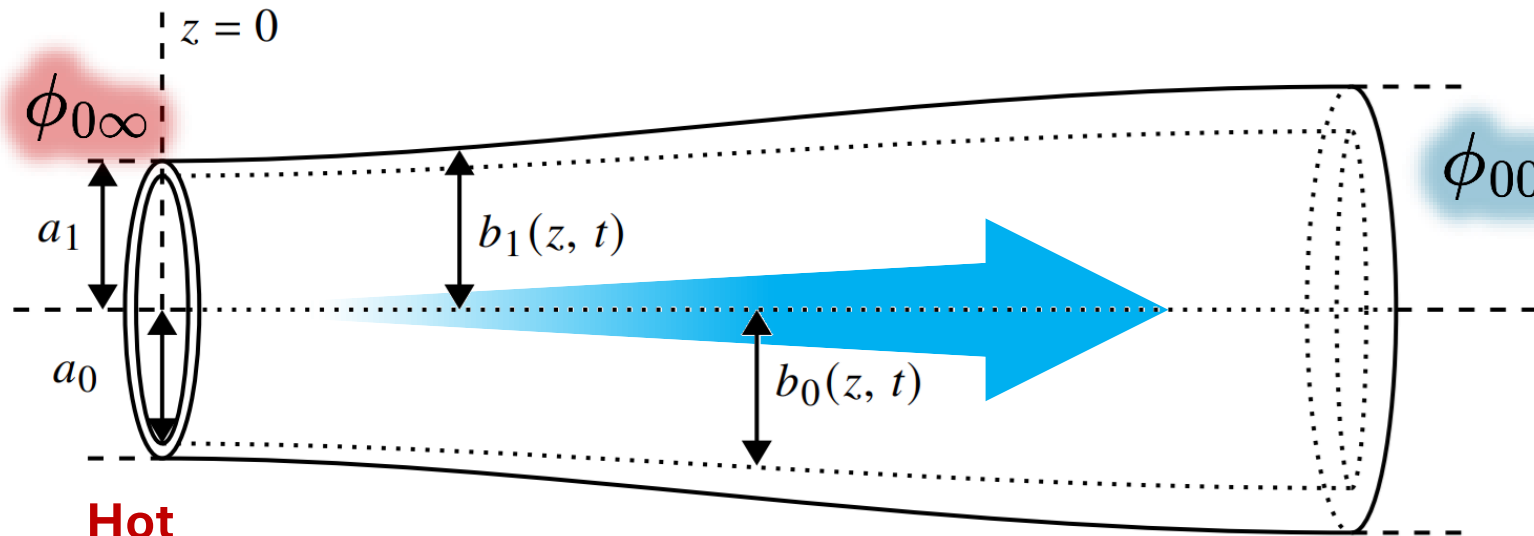
$$Le = \frac{\tau \kappa}{L^2} = \frac{\mu_l \kappa}{k \hat{\Pi}}$$

Typically  $\sim 10$ , so diffusion of heat dominates over compositional diffusion

$$\mathcal{D} = \frac{\hat{\Pi}}{c \Delta T}$$

Typically  $\sim 1/\Delta T$ , so importance of swelling and interstitial flows depends on magnitude of temperature differential

# Application: tubes of responsive gel



**Hot**

Shrunken tube, narrow lumen

$$T = T_C + \Delta T$$

**Cold**

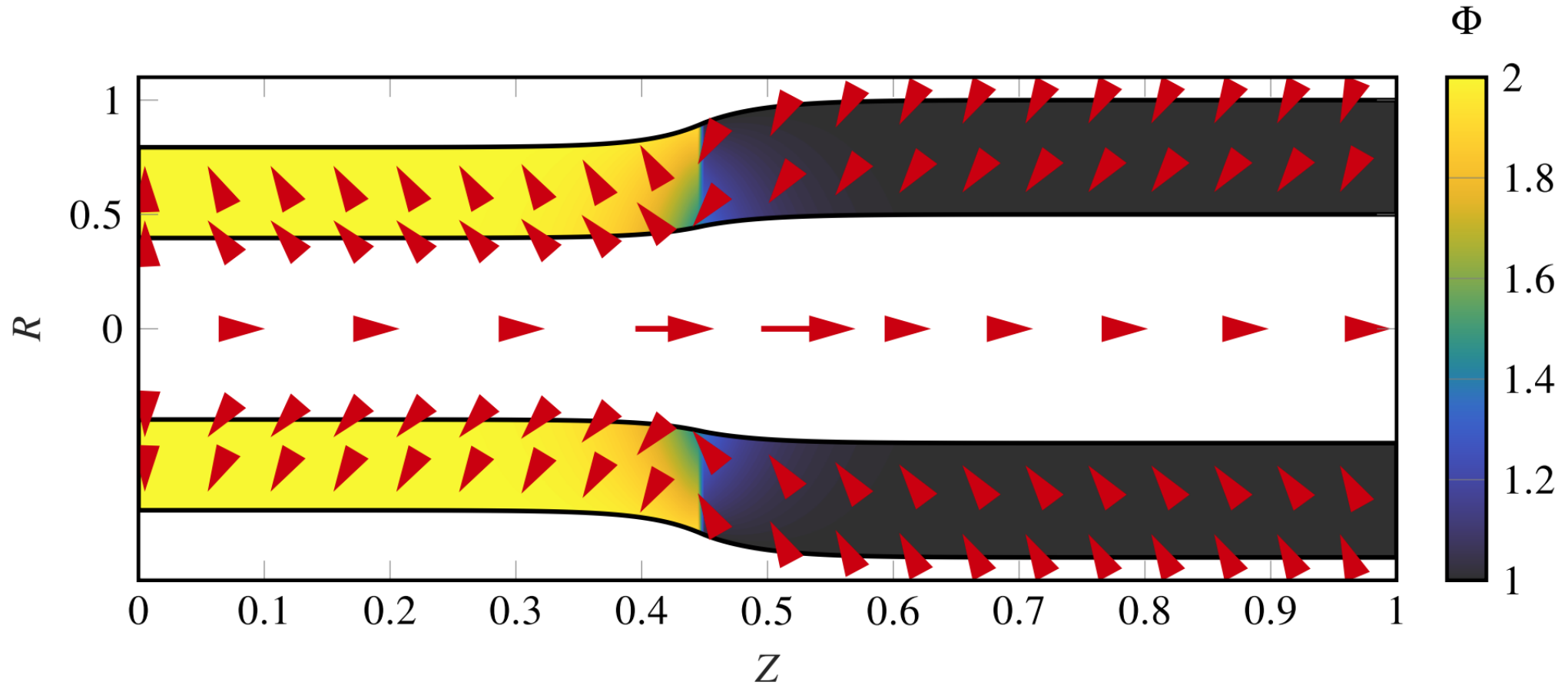
Relaxed tube, wide lumen

$$T \rightarrow T_C$$

$$T - T_C = \Delta T \left[ 2 \operatorname{erfc} \left( \frac{z}{2\sqrt{\kappa t}} \right) - 1 \right]$$

$$Z_C = 2 \operatorname{erfc}^{-1} \left( \frac{1}{2} \right) \sqrt{\kappa t}$$

# Flows through the tube





# Summary

- Modelling thermo-responsive gels requires an understanding of how the affinity of polymer chains for water changes with temperature
- This can be gained from microscopic models of the gel, or phenomenologically: spotting the LCST and how the equilibrium polymer fraction changes around it is enough to model the behaviour of these gels accurately
- We must be careful when there are sharp drying fronts, however – LENS can't handle these
- LENS does, however, allow us to carefully interpret all the terms in the heat transfer equation, closing the system

# With thanks to



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